

Ch. X: Supplement 2

- An alternative way in getting Eq. (2) on p. X-④.

$$-\frac{\hbar^2}{2m} \nabla^2 \psi + V(\vec{r}) \psi = E \psi$$

Bloch's theorem: $\psi_{\vec{k}}(\vec{r}) = e^{i\vec{k} \cdot \vec{r}} u_{\vec{k}}(\vec{r})$ with $u_{\vec{k}}(\vec{r})$

Eq. of $u_{\vec{k}}(\vec{r})$: (See Ch. IX)

$$= u_{\vec{k}}(\vec{r} + \vec{R})$$

$$-\frac{\hbar^2}{2m} (\nabla^2 + 2i\vec{k} \cdot \nabla - k^2) u_{\vec{k}}(\vec{r}) + V(\vec{r}) u_{\vec{k}}(\vec{r}) = E(\vec{k}) u_{\vec{k}}(\vec{r}) \quad (a)$$

This is of the form

$$\hat{H}_{\vec{k}} u_{\vec{k}}(\vec{r}) = E(\vec{k}) u_{\vec{k}}(\vec{r})$$

From Ch. IV (reciprocal lattice):

$$V(\vec{r}) = \sum_{\vec{G}_1} V(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}} \quad (\because V(\vec{r}) = V(\vec{r} + \vec{R}))$$

$$V(\vec{G}_1) = \frac{1}{\Omega_c} \int_{\Omega_c} V(\vec{r}) e^{-i\vec{G}_1 \cdot \vec{r}} d^3 r$$

Now, we turn (a) into a matrix problem.

② Expand $U_{\vec{k}}(\vec{r})$ in terms of a complete set.

Since $U_{\vec{k}}(\vec{r})$ is periodic, there is a natural choice of the set of functions.

$$U_{\vec{k}}(\vec{r}) = U_{\vec{k}}(\vec{r} + \vec{R}) \leftarrow \text{so as to satisfy Bloch's theorem}$$

One can write:

$$U_{\vec{k}}(\vec{r}) = \sum_{\vec{G}_1} U_{\vec{k}}(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}}$$

unknowns unknowns known set of functions
 $\{e^{i\vec{G}_1 \cdot \vec{r}}\}$ infinite many \vec{G}_1 's

1st term of (a):

$$\sum_{\vec{G}_1} -\frac{\hbar^2}{2m} \left(\nabla^2 + 2i\vec{k} \cdot \nabla - k^2 \right) U_{\vec{k}}(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}}$$

$$= \sum_{\vec{G}_1} -\frac{\hbar^2}{2m} \left(-\vec{G}_1^2 - 2\vec{k} \cdot \vec{G}_1 - k^2 \right) U_{\vec{k}}(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}}$$

$$= \sum_{\vec{G}_1} \frac{\hbar^2}{2m} |\vec{k} + \vec{G}_1|^2 U_{\vec{k}}(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}}$$

Eq.(a) becomes:

$$\sum_{\vec{G}_1} \frac{\hbar^2}{2m} |\vec{k} + \vec{G}_1|^2 U_{\vec{k}}(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}} + \sum_{\vec{G}_1} V(\vec{r}) U_{\vec{k}}(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}} = \sum_{\vec{G}_1} E(\vec{k}) U_{\vec{k}}(\vec{G}_1) e^{i\vec{G}_1 \cdot \vec{r}}$$

- Multiply by $e^{-i\vec{G}_1 \cdot \vec{r}}$

- Integrate $\int_{Q_c} d^3 r$ over a unit cell

$$\sum_{\vec{G}_1} \frac{\hbar^2}{2m} |\vec{k} + \vec{G}_1|^2 u_{\vec{k}}(\vec{G}_1) \frac{1}{V_e} \int_{V_e} e^{i(\vec{k} - \vec{G}_1) \cdot \vec{r}} d^3 r$$

$$+ \sum_{\vec{G}_1} u_{\vec{k}}(\vec{G}_1) \int_{V_e} V(\vec{r}) e^{-i(\vec{k}' - \vec{G}_1) \cdot \vec{r}} d^3 r = \sum_{\vec{G}_1} \epsilon(\vec{k}') u_{\vec{k}}(\vec{G}_1) \frac{1}{V_e} \int_{V_e} e^{i(\vec{k} - \vec{G}_1') \cdot \vec{r}} d^3 r$$

$$\Rightarrow \boxed{\frac{\hbar^2}{2m} |\vec{k} + \vec{G}_1'|^2 u_{\vec{k}}(\vec{G}_1') + \sum_{\vec{G}_1} V(\vec{G}_1' - \vec{G}_1) u_{\vec{k}}(\vec{G}_1) = \epsilon(\vec{k}) u_{\vec{k}}(\vec{G}_1')} \quad (b)$$

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$$\sum_{\vec{G}_1} \frac{\hbar^2}{2m} |\vec{k} + \vec{G}_1|^2 u_{\vec{k}}(\vec{G}_1) \delta_{\vec{G}_1, \vec{G}_1'} + \sum_{\vec{G}_1} V(\vec{G}_1' - \vec{G}_1) u_{\vec{k}}(\vec{G}_1) = \epsilon(\vec{k}) u_{\vec{k}}(\vec{G}_1')$$

$$\Rightarrow \sum_{\vec{G}_1} \left[\frac{\hbar^2}{2m} |\vec{k} + \vec{G}_1|^2 \delta_{\vec{G}_1, \vec{G}_1'} + V(\vec{G}_1' - \vec{G}_1) \right] u_{\vec{k}}(\vec{G}_1) = \epsilon(\vec{k}) u_{\vec{k}}(\vec{G}_1')$$

$$\Rightarrow \sum_{\vec{G}_1} \underbrace{H_{\vec{k}}(\vec{G}_1', \vec{G}_1)}_{\text{• } (\vec{G}_1', \vec{G}_1) \text{ matrix element}} u_{\vec{k}}(\vec{G}_1) = \epsilon(\vec{k}) u_{\vec{k}}(\vec{G}_1')$$

- (\vec{G}_1', \vec{G}_1) matrix element
- a matrix problem for each \vec{k}

$$\text{c.f. } \sum_i H_{ji} a_i = E a_j$$

* Look at diagonal elements of \hat{H}_k

i.e. $\hat{G}_1 = \hat{G}_1'$

a constant

$$\frac{\hbar^2}{2m} |\hat{k} + \hat{G}_1'|^2 + V(\vec{0}) = E(\hat{k} + \hat{G}_1') + V$$

* Off-diagonal elements $\hat{G}_1 \neq \hat{G}_1'$

$$V(\hat{G}_1' - \hat{G}_1)$$

$$\begin{array}{ccccccc} & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ - & \frac{1}{E^0(\hat{k}) + V} & e^{i\hat{G}_1 \cdot \vec{r}} & e^{i\hat{G}_2 \cdot \vec{r}} & e^{i\hat{G}_3 \cdot \vec{r}} & \dots & \dots \\ & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ e^{i\hat{G}_1 \cdot \vec{r}} & : & V(\hat{G}_1) & E^0(\hat{k} + \hat{G}_1) + V & V(\hat{G}_1 - \hat{G}_2) & V(\hat{G}_1 - \hat{G}_3) & \dots \\ & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ e^{i\hat{G}_2 \cdot \vec{r}} & : & V(\hat{G}_2) & V(\hat{G}_2 - \hat{G}_1) & E^0(\hat{k} + \hat{G}_2) + V & V(\hat{G}_2 - \hat{G}_3) & \dots \\ & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ e^{i\hat{G}_3 \cdot \vec{r}} & : & V(\hat{G}_3) & V(\hat{G}_3 - \hat{G}_1) & V(\hat{G}_3 - \hat{G}_2) & E^0(\hat{k} + \hat{G}_3) + V & \dots \\ & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \end{array}$$

$$\left(\begin{array}{c} \vdots \\ \vdots \\ \vdots \\ \vdots \\ \vdots \end{array} \right) \left(\begin{array}{c} U_{\hat{k}}(\vec{0}) \\ U_{\hat{k}}(\hat{G}_1) \\ U_{\hat{k}}(\hat{G}_2) \\ U_{\hat{k}}(\hat{G}_3) \\ \vdots \end{array} \right) = E(\hat{k}) \left(\begin{array}{c} U_{\hat{k}}(\vec{0}) \\ U_{\hat{k}}(\hat{G}_1) \\ U_{\hat{k}}(\hat{G}_2) \\ U_{\hat{k}}(\hat{G}_3) \\ \vdots \end{array} \right) \quad (C)$$

→ Same as Eq.(b)

Start again from Eq.(b):

$$\frac{\hbar^2}{2m} |\vec{k} + \vec{G}'|^2 u_{\vec{k}}(\vec{G}') + V(\vec{0}) u_{\vec{k}}(\vec{G}') + \sum_{\vec{G} \neq \vec{G}'} V(\vec{G}' - \vec{G}) u_{\vec{k}}(\vec{G}) = E(\vec{k}) u_{\vec{k}}(\vec{G})$$

$$\Rightarrow \left(E(\vec{k}) - V(\vec{0}) - \frac{\hbar^2}{2m} |\vec{k} + \vec{G}'|^2 \right) u_{\vec{k}}(\vec{G}') = \sum_{\substack{\vec{G} \\ (\vec{G} \neq \vec{G}')}} V(\vec{G}' - \vec{G}) u_{\vec{k}}(\vec{G})$$

$$= \sum_{\substack{\vec{G}'' \\ (\vec{G}'' \neq 0)}} V(\vec{G}'') u_{\vec{k}}(\vec{G}' - \vec{G}'') \quad (d)$$

which is the equation (Eq.(2)) on page IX-4.